

## Efficient Numerical Schemes for Computations of European Options with Transaction Costs

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**ABSTRACT.** This paper aims to find numerical solutions of the non-linear Black-Scholes partial differential equation (PDE), which often appears in financial markets, for European option pricing in the appearance of the transaction costs. Here we exploit the transformations for the computational purpose of a non-linear Black-Scholes PDE to modify as a non-linear parabolic type PDE with reliable initial and boundary conditions for call and put options. Several schemes are derived rigorously using the finite volume method (FVM) and finite difference method (FDM), which is the novelty of this paper. Stability and consistency analysis assure the convergence of these schemes. We apply these schemes to various volatility models, such as the Leland, Boyle and Vorst, Barles and Soner, and Risk-adjusted pricing methodology (RAPM). All the schemes are tested numerically. The convergence of the obtained results is observed, and we find that they are also reliable. Finally, we display all the approximate results together with the exact values through graphical and tabular representations.

### 1. INTRODUCTION

Understanding and accurately evaluating transaction costs in a financial market is vital for security trading, asset pricing, stock market regulation, and many other issues. During the last few decades, pricing options more accurately after including realistic assumptions—such as transaction cost, getting more importance from both the traders and the investors.

The literature's [1–6], contains descriptive discussions of options. Fischer Black and Myron Scholes [7] worked jointly, and first disclosed the concept of the Black-Scholes model for *options pricing and corporate liabilities*, and was published in 1973, while Robert Merton [8] advanced this model in the article "Theory of rational option pricing" in the same year. Their derived equation is based on the assumption that there are no fees for buying and selling options and stocks, as well as no trade barriers (i.e., no commissions and transaction costs). In other words, this model makes

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a friction-less assumption (which is indispensable, as actual costs correlated with practical market applications) to implement a hedging plan for any contingent claim of the European type.

Various studies have been conducted about the linear Black-Scholes model [9–15] though it adopts the unrealistic assumption of no transaction costs. Several studies have been attempted to evaluate the price of European options [16–23], American options [24–28], Asian options [29, 30], and Barrier options [31] in a completely friction-less market. Recently, the fractional Black-Scholes model [32–34] received some attention.

Contrastingly, the non-linear Black-Scholes PDE, where the non-linear term denotes the presence of transaction costs, is of great importance to our contemporary world over some time both in terms of approach and applicability. Several models [35] consider transaction costs: Leland model, Paras, and Avellaneda model, Boyle and Vorst model, Hodges and Neuberger model, Barles and Soner model, and RAPM (Risk-adjusted pricing methodology) model. If the transaction cost parameters are equal to zero, all of these non-linear transaction cost models are unvarying with the linear model.

Soner et al. [36] showed that there is no nontrivial hedging portfolio for option pricing with transaction costs. They also suggested that the best hedging strategy is buying an asset and taking on it for a certain period as a call or put option. Leland [37] inaugurates the idea of using transaction costs at discrete times. He also indicated that the hedging error could be minimized if the length of re-balancing frequency approaches zero. Later, Boyle and Vorst [38] demonstrated further in a discrete-time framework with a binomial tree model for the option prices with proportional transaction costs, and it is pretty accurate for possible parameter values. Besides, Dewynne et al. [39] considered path-dependent and exotic options with transaction costs. Recently, asymptotic analysis [40] and Markov chain approximation [41] were also studied for pricing European options with transaction costs in some previous literature.

On the other hand, few researchers [42–46, 49–51] paid their attention to solve the non-linear Black-Scholes equation numerically. For example, the exponential time differencing (ETD) method [44] was applied to solve the non-linear Black-Scholes model for pricing American options with a highly stable and efficient transaction cost. Lesmana and Wang [45] developed the numerical method based on an upwind finite difference scheme for a non-linear parabolic PDE, and they attempted to pricing European options under transaction costs. Ankudinova and Ehrhardt [46] focused on the non-linear Black-Scholes equation for European call options using several transaction cost models as well as Crank–Nicolson and Rigal compact schemes. R. L. Valkov [47] has solved the non-linear Black-Scholes-Bellman model numerically as well as discuss the monotonicity and consistency of his suggested scheme in considerable detail. A monotone finite volume spatial discretization and a second-order predictor-corrector scheme in time are considered by Radoslas Valkov [48] to handle the Black-Scholes equation with uncertain volatility and dividend. The applicability of implicit

numerical schemes for the valuation of contingent claims in non-linear Black–Scholes models has been discussed by Pascal Heider [49]. He also studied the practical implications of the derived stability criteria on relevant numerical examples. He claimed that if certain stability requirements are satisfied, it is possible to construct convergent implicit algorithms for non-linear Black–Scholes equations. Ekaterina Dremkova and Matthias Ehrhardt [50] have solved non-linear Black–Scholes equations for American options with a non-linear volatility function using various compact finite difference techniques to improve the order of the accuracy. The existence and uniqueness of solutions to the well-known non-linear Black–Scholes equation have been demonstrated by Naoyuki Ishimura [51] for both in the classical and weak senses.

However, in this paper, we work on approximating non-linear Black–Scholes PDE for valuing European options when there are transaction costs. For this, we organize the present research work as follows: we modify the original model into parabolic type PDE exploiting the transformations [46] which are written in section 2. A brief description of different volatility models is given in section 3 subsequently. Section 4 is devoted to discretize the transformed parabolic type equation by using some numerical schemes. Stability and consistency analysis are included in sections 5 and 6, respectively. In section 7, numerical examples are given to show the efficacy of the proposed schemes. Subsequently, a general conclusion is drawn in section 8. Finally, all relevant references are included.

## 2. THE MODEL EQUATION

This section considers a non-linear Black–Scholes PDE and modifies it to a non-linear parabolic type equation with appropriate and available transformations, which would be easy to compute numerically. Let us consider the non-linear Black–Scholes PDE [46],

$$\frac{\partial F}{\partial t} + rS \frac{\partial F}{\partial S} + \frac{1}{2} \tilde{\sigma}^2 S^2 \frac{\partial^2 F}{\partial S^2} - rF = 0, 0 < S < \infty, t \in (0, T) \quad (1)$$

subject to the terminal and boundary conditions for European call and put options:

$F(S, T) = \max(S - K, 0)$ ,  $F(S, t) = 0$  when  $S = 0$ ,  $F(S, t) = S - Ke^{-r(T-t)}$ , when  $S \rightarrow \infty$  and  $F(S, T) = \max(K - S, 0)$ ,  $F(S, t) = Ke^{-r(T-t)}$ , when  $S = 0$ ,  $F(S, t) = 0$ , when  $S \rightarrow \infty$  respectively. Throughout this paper, we use the notations:  $F = F(S, t)$  = the option price,  $S$  = stock price,  $K$  = strike price,  $T$  = maturity time,  $r$  = interest rate,  $t$  = time in years, and  $\tilde{\sigma} = \tilde{\sigma} \left( t, S, \frac{\partial F}{\partial S}, \frac{\partial^2 F}{\partial S^2} \right)$  depends on the volatility model.

Now consider the transformations [46] as given below,

$$y = \ln(K^{-1}S), \tau = \frac{1}{2}\sigma^2(T-t) \text{ and } u(y, t) = K^{-1}e^{-y}F(S, t)$$

and substituting these into Equation (1) to obtain the following non-linear parabolic PDE

$$\frac{\partial u}{\partial \tau} = \frac{2r}{\sigma^2} \frac{\partial u}{\partial y} + \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \left( \frac{\partial^2 u}{\partial y^2} + \frac{\partial u}{\partial y} \right), y_{\min} < y < y_{\max}, \tau \in \left( 0, \frac{\sigma^2}{2} T \right) \quad (2)$$

with the modified initial and boundary conditions for European call and put options:

$u(y, 0) = \max(1 - e^{-y}, 0)$  as  $y \in (-\infty, \infty)$ ,  $u(y, \tau) = 0$  as  $y \rightarrow -\infty$ ,  $u(y, \tau) = 1 - e^{-(y+2r\tau/\sigma^2)}$  as  $y \rightarrow \infty$ , and  $u(y, 0) = \max(e^{-y} - 1, 0)$  as  $y \in (-\infty, \infty)$ ,  $u(y, \tau) = e^{-(y+2r\tau/\sigma^2)}$  as  $y \rightarrow -\infty$ ,  $u(y, \tau) = 0$  as  $y \rightarrow \infty$ , respectively.

### 3. VOLATILITY MODELS

This section concerns four stochastic volatility models to discretize the non-linear Black-Scholes PDE, whose solution provides the option price for transaction fees. We give a short description, but details are available in some previous literature [46].

**Leland Volatility Model (LVM).** Leland [37] developed a technique for replicating options in the presence of transactions costs for a small time interval. He proposed that the option price is the solution of the non-linear Black-Scholes Equation (1) but with the adjusted volatility [46] as follows:

$$\tilde{\sigma} = \sigma \sqrt{1 + \sqrt{\frac{2}{\pi}} \frac{\mu}{\sigma \sqrt{\Delta t}} \text{sign}(F_{SS})} \quad (3)$$

where,  $\sigma$  is the original volatility,  $\mu$  is the round-trip transaction cost per unit dollar of the transaction, and  $\Delta t$  is the transaction frequency. In this formula, both  $\mu$  and  $\Delta t$  are assumed to be small while keeping the ratio  $\frac{\mu}{\sqrt{\Delta t}}$  of order one.

**Boyle and Vorst Volatility Model (BVVM).** Boyle and Vorst [38] derived a method for calculating option prices in a discrete-time where option price meets to Black-Scholes price with the modified volatility [46] given by

$$\tilde{\sigma} = \sigma \sqrt{1 + \frac{\mu}{\sigma \sqrt{\Delta t}} \text{sign}(F_{SS})} \quad (4)$$

where,  $\sigma$ ,  $\mu$ , and  $\Delta t$  represents the same meaning as Leland.

**Barles and Soner Volatility Model (BSVM).** Barles and Soner [43] evolved a model using the utility function approach of Hodges and Neuberger [52] along with asymptotic analysis of partial differential equations. For this case, the formula for the modified volatility [46] is given by

$$\tilde{\sigma} = \sigma \sqrt{1 + e^{r(T-t)} a^2 S^2 F_{SS}} \quad (5)$$

where,  $\mu = a\sqrt{\epsilon}$  is the round-trip transaction cost per unit dollar of the transaction for some constant  $a > 0$  and  $\epsilon \rightarrow 0$ .

**RAPM Volatility Model (RAPMVM).** Kratka [53] took the first step for this model and later improved by Jandačka and Ševčovič [54]. Here the modified volatility is of the form [46]

$$\tilde{\sigma} = \sigma \sqrt{1 + 3 \times \sqrt[3]{\frac{C^2 M}{2\pi} S F_{SS}}} \quad (6)$$

where,  $M \geq 0$  is the transaction cost measure and  $C \geq 0$  is the risk premium measure.

#### 4. DERIVATIONS OF COMPUTATIONAL SCHEMES

In this section, we derive five computational schemes, in detail, for Equation (2) using two well-known numerical methods.

**4.1. Dufort-Frankel Finite Difference Scheme.** The Dufort-Frankel FD scheme [55] can be applied to solve various kinds of problems which occur in finance. This scheme is a multi-step method, and requires another scheme for simulating the first temporal vector. In this formulation  $\frac{\partial u}{\partial \tau}$ ,  $\frac{\partial u}{\partial y}$ , and  $\frac{\partial^2 u}{\partial y^2}$  are discretized by central difference and  $u_i^j$  is replaced by  $(u_i^{j+1} + u_i^{j-1})/2$ . Thus, discretizing Equation (2) by Dufort-Frankel FDM, we obtain

$$\begin{aligned} \frac{1}{2\Delta\tau} (u_i^{j+1} - u_i^{j-1}) &= \left(\frac{\tilde{\sigma}}{\sigma}\right)^2 \left[ \frac{1}{(\Delta y)^2} (u_{i-1}^j - (u_i^{j+1} + u_i^{j-1}) + u_{i+1}^j) \right] \\ &+ \left(\frac{\tilde{\sigma}}{\sigma}\right)^2 \left[ \frac{1}{2\Delta y} (u_{i+1}^j - u_{i-1}^j) \right] + \frac{r}{\sigma^2 \Delta y} (u_{i+1}^j - u_{i-1}^j) \end{aligned}$$

or, equivalently

$$\begin{aligned} u_i^{j+1} &= u_i^{j-1} + \frac{2r(\Delta\tau)(\Delta y)}{\sigma^2} (u_{i+1}^j - u_{i-1}^j) + \frac{\Delta\tau}{\Delta y} \left(\frac{\tilde{\sigma}}{\sigma}\right)^2 (u_{i+1}^j - u_{i-1}^j) \\ &+ \frac{2(\Delta\tau)}{(\Delta y)^2} \left(\frac{\tilde{\sigma}}{\sigma}\right)^2 (u_{i-1}^j - u_i^{j+1} - u_i^{j-1} + u_{i+1}^j) \end{aligned}$$

which can be written as

$$u_i^{j+1} = a_i u_{i-1}^j + b_i u_{i+1}^j + c_i u_i^{j-1}; i = 0, 1, 2, \dots, n-1; j = 0, 1, 2, \dots, m-1 \quad (7)$$

where

$$\begin{aligned} a_i &= \left[ (\Delta y)^2 + 2(\Delta\tau) \left(\frac{\tilde{\sigma}}{\sigma}\right)^2 \right]^{-1} \times \left[ (\Delta\tau) \left(\frac{\tilde{\sigma}}{\sigma}\right)^2 (2 - \Delta y) - (\Delta y)(\Delta\tau) \frac{2r}{\sigma^2} \right], \\ b_i &= \left[ (\Delta y)^2 + 2(\Delta\tau) \left(\frac{\tilde{\sigma}}{\sigma}\right)^2 \right]^{-1} \times \left[ (\Delta\tau) \left(\frac{\tilde{\sigma}}{\sigma}\right)^2 (2 + \Delta y) + (\Delta y)(\Delta\tau) \frac{2r}{\sigma^2} \right], \end{aligned}$$

and

$$c_i = \left[ (\Delta y)^2 + 2(\Delta\tau) \left(\frac{\tilde{\sigma}}{\sigma}\right)^2 \right]^{-1} \times \left[ (\Delta y)^2 - 2(\Delta\tau) \left(\frac{\tilde{\sigma}}{\sigma}\right)^2 \right]$$

which is our proposed Dufort-Frankel Finite Difference Scheme (DFFDS).

**4.2. Laasonen Finite Difference Scheme.** The Laasonen finite difference scheme [55] can be applied to solve linear and non-linear partial differential equations. This method metamorphosed partial differential equations into a system of linear algebraic equations. In this formulation  $\frac{\partial u}{\partial \tau}$  is approximated by a central differencing at a step  $\frac{\Delta \tau}{2}$ , and  $\frac{\partial u}{\partial y}, \frac{\partial^2 u}{\partial y^2}$  are approximated by central differences at time levels  $j + 1$ . Now the discretized form of Equation (2) is as follows

$$\frac{1}{\Delta \tau} \left( u_i^{j+1} - u_i^j \right) = \frac{1}{2(\Delta y)^2} \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \left[ 2 \left( u_{i-1}^{j+1} - 2u_i^{j+1} + u_{i+1}^{j+1} \right) + \Delta y \left( u_{i+1}^{j+1} - u_{i-1}^{j+1} \right) \right] + \frac{r}{\sigma^2 \Delta y} \left( u_{i+1}^{j+1} - u_{i-1}^{j+1} \right)$$

After simplification, we get

$$\left[ \frac{r\Delta\tau}{\Delta y\sigma^2} + \frac{\Delta\tau}{2(\Delta y)^2} \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 (\Delta y - 2) \right] u_{i-1}^{j+1} + \left[ 1 + \frac{2\Delta\tau}{(\Delta y)^2} \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \right] u_i^{j+1} - \left[ \frac{r\Delta\tau}{\Delta y\sigma^2} + \frac{\Delta\tau}{2(\Delta y)^2} \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 (\Delta y + 2) \right] u_{i+1}^{j+1} = u_i^j$$

The above equation reduces to

$$d_i u_{i-1}^{j+1} + (1 + e_i) u_i^{j+1} + f_i u_{i+1}^{j+1} = u_i^j; i = 0, 1, 2, \dots, n-1; j = 0, 1, 2, \dots, m-1 \quad (8)$$

where

$$d_i = \frac{r\Delta\tau}{\Delta y\sigma^2} + \frac{\Delta\tau}{2(\Delta y)^2} \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 (\Delta y - 2), e_i = \frac{2\Delta\tau}{(\Delta y)^2} \left( \frac{\tilde{\sigma}}{\sigma} \right)^2$$

and

$$f_i = -\frac{r\Delta\tau}{\Delta y\sigma^2} - \frac{\Delta\tau}{2(\Delta y)^2} \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 (\Delta y + 2)$$

**4.3. Finite Volume Schemes.** The finite volume scheme is a scheme of solving different kinds of time-dependent or independent partial differential equations in algebraic equations. In this scheme, we divide the physical space into a finite number of control volumes. In this section, we describe it in a few lines, but details are available in the previous study [56] conducted by Malalasekera et al.

Applying the finite volume integration in Equation (2) over a control volume (CV) with a finite time step  $\Delta \tau$ , we obtain

$$\int_{\tau}^{\tau+\Delta\tau} \int_{CV} \frac{\partial u}{\partial \tau} dV d\tau = \left( \frac{2r}{\sigma^2} + \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \right) \int_{\tau}^{\tau+\Delta\tau} \int_{CV} \frac{\partial u}{\partial y} dV d\tau + \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \int_{\tau}^{\tau+\Delta\tau} \int_{CV} \frac{\partial^2 u}{\partial y^2} dV d\tau$$

After rearranging, we get

$$\int_{CV} \left[ \int_{\tau}^{\tau+\Delta\tau} \frac{\partial u}{\partial \tau} d\tau \right] dV = \left( \frac{2r}{\sigma^2} + \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \right) \int_{\tau}^{\tau+\Delta\tau} \left[ \int_{CV} \frac{\partial u}{\partial y} dV \right] d\tau + \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \int_{\tau}^{\tau+\Delta\tau} \left[ \int_{CV} \frac{\partial^2 u}{\partial y^2} dV \right] d\tau$$

Applying Gauss's divergence theorem, the above equation leads

$$(u_P - u_P^0) \Delta V = \frac{1}{2} \left( \frac{2r}{\sigma^2} + \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \right) A \int_{\tau}^{\tau+\Delta\tau} (u_E - u_W) d\tau + \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \int_{\tau}^{\tau+\Delta\tau} \left[ \left( A \frac{u_E - u_P}{\delta y_{PE}} \right) - \left[ \left( A \frac{u_P - u_W}{\delta y_{WP}} \right) \right] \right] d\tau \quad (9)$$

For  $0 \leq \theta \leq 1$ , we assume

$$\int_{\tau}^{\tau+\Delta\tau} u_P d\tau = [\theta u_P + (1 - \theta) u_P^0] \Delta\tau \quad (10)$$

Applying Equation (10) into Equation (9) and dividing by we get

$$(u_P - u_P^0) \frac{\Delta y}{\Delta\tau} = \frac{1}{2} \left( \frac{2r}{\sigma^2} + \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \right) [\theta (u_E - u_W) + (1 - \theta) (u_E^0 - u_W^0)] + \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \theta \left( \frac{u_E - u_P}{\delta y_{PE}} - \frac{u_P - u_W}{\delta y_{WP}} \right) + \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 (1 - \theta) \left( \frac{u_E^0 - u_P^0}{\delta y_{PE}} - \frac{u_P^0 - u_W^0}{\delta y_{WP}} \right) \quad (11)$$

For convenience, we put  $\delta y_{WP} = \delta y_{PE} = \Delta y$  on the following three schemes.

*Explicit Scheme.* Substitution of  $\theta = 0$  into Equation (11) gives the following explicit discretized equation,

$$(u_P - u_P^0) \frac{\Delta y}{\Delta\tau} = \frac{1}{2} \left( \frac{2r}{\sigma^2} + \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \right) (u_E^0 - u_W^0) + \frac{1}{\Delta y} \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 (u_E^0 - 2u_P^0 + u_W^0)$$

This equation may be re-writtens as

$$u_P = \alpha_i u_W^0 + (1 + \beta_i) u_P^0 + \gamma_i u_E^0 \quad (12)$$

where

$$\alpha_i = \frac{\Delta\tau}{(\Delta y)^2} \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 - \frac{\Delta\tau}{2\Delta y} \left( \frac{2r}{\sigma^2} + \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \right), \quad \beta_i = -\frac{2\Delta\tau}{(\Delta y)^2} \left( \frac{\tilde{\sigma}}{\sigma} \right)^2$$

and

$$\gamma_i = \frac{\Delta\tau}{(\Delta y)^2} \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 + \frac{\Delta\tau}{2\Delta y} \left( \frac{2r}{\sigma^2} + \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \right)$$

which is the desired Finite Volume Explicit Scheme (FVES).

*Crank-Nicolson Scheme.* Putting  $\theta = \frac{1}{2}$  into Equation (11), we get the following Crank-Nicolson discretized equation,

$$\begin{aligned} (u_P - u_P^0) \frac{\Delta y}{\Delta \tau} = & \frac{1}{4} \left( \frac{2r}{\sigma^2} + \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \right) (u_E - u_W + u_E^0 - u_W^0) \\ & + \frac{1}{2\Delta y} \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 (u_E - 2u_P + u_W + u_E^0 - 2u_P^0 + u_W^0) \end{aligned}$$

After simplification, we get the following equation

$$\lambda_i u_W + (1 + \xi_i) u_P + \eta_i u_E = -\lambda_i u_P^0 + (1 - \xi_i) u_P^0 - \eta_i u_E^0 \quad (13)$$

where

$$\lambda_i = \frac{\Delta \tau}{4\Delta y} \left( \frac{2r}{\sigma^2} + \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \right) - \frac{\Delta \tau}{2(\Delta y)^2} \left( \frac{\tilde{\sigma}}{\sigma} \right)^2, \xi_i = \frac{\Delta \tau}{(\Delta y)^2} \left( \frac{\tilde{\sigma}}{\sigma} \right)^2$$

and

$$\eta_i = - \left[ \frac{\Delta \tau}{4\Delta y} \left( \frac{2r}{\sigma^2} + \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \right) + \frac{\Delta \tau}{2(\Delta y)^2} \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \right]$$

which is the proposed Finite Volume Crank-Nicolson Scheme (FVCNS).

*Fully Implicit Scheme.* Substitution of  $\theta = 1$  into Equation (11) leads to the following form:

$$(u_P - u_P^0) \frac{\Delta y}{\Delta \tau} = \frac{1}{2} \left( \frac{2r}{\sigma^2} + \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \right) (u_E - u_W) + \frac{1}{\Delta y} \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 (u_E - 2u_P + u_W)$$

and the reduced formula is then

$$q_i u_W + (1 + r_i) u_P + s_i u_E = u_P^0 \quad (14)$$

where

$$q_i = \frac{\Delta \tau}{2\Delta y} \left( \frac{2r}{\sigma^2} + \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \right) - \frac{\Delta \tau}{(\Delta y)^2} \left( \frac{\tilde{\sigma}}{\sigma} \right)^2, r_i = \frac{2\Delta \tau}{(\Delta y)^2} \left( \frac{\tilde{\sigma}}{\sigma} \right)^2$$

and

$$s_i = - \frac{\Delta \tau}{2\Delta y} \left( \frac{2r}{\sigma^2} + \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \right) - \frac{\Delta \tau}{(\Delta y)^2} \left( \frac{\tilde{\sigma}}{\sigma} \right)^2$$

which is our proposed Finite Volume Fully Implicit Scheme (FVFIS).

## 5. STABILITY OF THE NUMERICAL SCHEMES

To test the stability of the derived schemes in section 4, with the help of the Von-Neumann stability method [55], let us consider a Fourier component for  $u_i^j$  and  $u_P^0$  as

$$u_i^j = U^j e^{l\theta i} \text{ and } u_P^0 = U^j e^{l\theta i} \quad (15)$$

where  $l = \sqrt{-1}$ , i.e., imaginary unit,  $U^j$  is the amplitude at a time level  $j$ ,  $\theta (= R\Delta y)$  is the phase angle,  $R$  is the wave number in the  $x$ -direction, and  $i$  represents the index of the node.

Similarly,

$$\begin{aligned} u_{i\pm 1}^{j\mp 1} &= U^{j\mp 1} e^{l\theta(i\pm 1)} & u_W^0 &= U^j e^{l\theta(i-1)} & u_E^0 &= U^j e^{l\theta(i+1)} & u_P &= U^{j+1} e^{l\theta i} \\ u_W &= U^{j+1} e^{l\theta(i-1)} & u_E &= U^{j+1} e^{l\theta(i+1)} \end{aligned} \quad (16)$$

For convenience, let us suppose that  $G = \frac{U^{j+1}}{U^j}$ . Thus, the stability requirement is  $|G|^2 \leq 1$ . Applying Equation (15) and Equation (16) into Equation (7), and dividing by  $e^{l\theta i}$ , we get

$$\begin{aligned} |G|^2 &= \frac{1}{4} \left[ \left\{ 4\Delta\tau \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \cos\theta \pm \sqrt{A} \right\}^2 + 4(\Delta\tau)^2 (\Delta y)^2 \left( \frac{2r}{\sigma^2} + \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \right)^2 (1 - \cos^2\theta) \right] \\ &\quad \times \left( (\Delta y)^2 + 2\Delta\tau \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \right)^{-2} \end{aligned} \quad (17)$$

where

$$\begin{aligned} A &= 16(\Delta\tau)^2 \left( \frac{\tilde{\sigma}}{\sigma} \right)^4 \cos^2\theta - 4(\Delta\tau)^2 (\Delta y)^2 \left( \frac{2r}{\sigma^2} + \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \right)^2 (1 - \cos^2\theta) + 4(\Delta y)^4 \\ &\quad - 16(\Delta\tau)^2 \left( \frac{\tilde{\sigma}}{\sigma} \right)^4 + 16(\Delta\tau)^2 (\Delta y) \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \left( \frac{2r}{\sigma^2} + \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \right) \cos\theta \sqrt{1 - \cos^2\theta} \end{aligned}$$

For extremum value of  $|G|^2$ , solving  $\frac{d|G|^2}{d(\cos\theta)} = 0$  for  $\cos\theta$ , and substituting it into  $\frac{d^2|G|^2}{d(\cos\theta)^2} < 0$ . Then from Equation (17), we cannot confirm that the maximum value of  $|G|^2$  would occur. However, the extreme values of  $\cos\theta$  must yet be investigated. For  $\cos\theta = 1$ , Equation (17) gives  $|G|^2 = 1$ , and the stability requirement is satisfied. For  $\cos\theta = -1$ , Equation (17) also yields  $|G|^2 = 1$  and, and the stability requirement is satisfied. Thus, the DFFDS proposed in Equation (7) is unconditionally stable.

Similarly, we can show that LFDS and FVFIS wrote in Equation (8) and Equation (14), respectively, both are unconditionally stable.

Again, applying Equation (15) and Equation (16) into Equation (12) and dividing by  $e^{l\theta i}$ , we get

$$G = 1 + \frac{2\Delta\tau}{(\Delta y)^2} \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 (\cos\theta - 1) + \frac{\Delta\tau}{\Delta y} \left( \frac{2r}{\sigma^2} + \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \right) \sin\theta$$

Then we may obtain easily,

$$|G|^2 = \left\{ 1 + \frac{2\Delta\tau}{(\Delta y)^2} \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 (\cos\theta - 1) \right\}^2 + \left( \frac{\Delta\tau}{\Delta y} \right)^2 \left( \frac{2r}{\sigma^2} + \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \right)^2 (1 - \cos^2\theta) \quad (18)$$

For extremum value of  $|G|^2$  such that  $\frac{d|G|^2}{d(\cos\theta)} = 0$ , we can find

$$\cos\theta = \frac{1}{\Delta\tau} \times \left[ 2 \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 - 4 \frac{\Delta\tau}{(\Delta y)^2} \left( \frac{\tilde{\sigma}}{\sigma} \right)^4 \right] \times \left( \left( \frac{2r}{\sigma^2} + \left( \frac{\tilde{\sigma}}{\sigma} \right)^2 \right)^2 - 4 \frac{\Delta\tau}{(\Delta y)^2} \left( \frac{\tilde{\sigma}}{\sigma} \right)^4 \right)^{-1} \quad (19)$$

Considering  $\frac{d^2|G|^2}{d(\cos\theta)^2} < 0$ , and substituting the value of  $\cos\theta$  from Equation (19) into Equation (18), which does not provide us the maximum value of  $|G|^2$ . But, the extreme values of  $\cos\theta$  must be investigated. For  $\cos\theta = 1$ , Equation (18) gives  $|G|^2 = 1$ , and the stability requirement is satisfied. For  $\cos\theta = -1$ , Equation (18) yields  $|G|^2 = \left\{1 - \frac{4\Delta\tau}{(\Delta y)^2} \left(\frac{\tilde{\sigma}}{\sigma}\right)^2\right\}^2$  and, imposing the requirement of  $|G|^2 \leq 1$ , yields, FVES in Equation (12) is conditionally stable and the condition is

$$\left(\frac{\tilde{\sigma}}{\sigma}\right)^2 \leq \frac{(\Delta y)^2}{2\Delta\tau} \quad (20)$$

Similarly, we can state that FVCNS, Equation (13) is also conditionally stable and the condition is

$$\left(\frac{\tilde{\sigma}}{\sigma}\right)^2 \leq \frac{(\Delta y)^2}{\Delta\tau} \quad (21)$$

## 6. CONSISTENCY OF THE NUMERICAL SCHEMES

For consistency, the finite difference equation (FDE) approximation of a PDE must reduce to the original PDE as the step sizes approach zero [55].

Now expanding each  $u(y, \tau)$  in a Taylor series expansion about  $u_i^j$ , we get

$$u_i^{j+1} = u_i^j + \Delta\tau \frac{\partial u}{\partial \tau} + \frac{(\Delta\tau)^2}{2!} \frac{\partial^2 u}{\partial \tau^2} + \frac{(\Delta\tau)^3}{3!} \frac{\partial^3 u}{\partial \tau^3} + O(\Delta\tau)^4 \quad (22)$$

$$\begin{aligned} u_{i+1}^{j+1} = & u_i^j + \Delta\tau \frac{\partial u}{\partial \tau} + \Delta y \frac{\partial u}{\partial y} + \frac{1}{2!} \left( \Delta\tau \frac{\partial}{\partial \tau} + \Delta y \frac{\partial}{\partial y} \right)^2 u \\ & + \frac{1}{3!} \left( \Delta\tau \frac{\partial}{\partial \tau} + \Delta y \frac{\partial}{\partial y} \right)^3 u + O[(\Delta\tau)^4, (\Delta y)^4] \end{aligned} \quad (23)$$

$$\begin{aligned} u_{i-1}^{j+1} = & u_i^j + \Delta\tau \frac{\partial u}{\partial \tau} - \Delta y \frac{\partial u}{\partial y} + \frac{1}{2!} \left( \Delta\tau \frac{\partial}{\partial \tau} - \Delta y \frac{\partial}{\partial y} \right)^2 u \\ & + \frac{1}{3!} \left( \Delta\tau \frac{\partial}{\partial \tau} - \Delta y \frac{\partial}{\partial y} \right)^3 u + O[(\Delta\tau)^4, (\Delta y)^4] \end{aligned} \quad (24)$$

Applying Equations (22), (23), and (24) into Equation (8) yields

$$\begin{aligned} & (d_i + e_i + f_i) u_i^j + (1 + d_i + e_i + f_i) \Delta\tau \frac{\partial u}{\partial \tau} + (1 + d_i + e_i + f_i) \frac{(\Delta\tau)^2}{2} \frac{\partial^2 u}{\partial \tau^2} \\ & + (-d_i + f_i) \Delta y \frac{\partial u}{\partial y} + (-d_i + f_i) \Delta\tau \Delta y \frac{\partial^2 u}{\partial \tau \partial y} + (d_i + f_i) \frac{(\Delta y)^2}{2} \frac{\partial^2 u}{\partial y^2} \\ & + O[(\Delta\tau)^3, (\Delta y)^3] = 0 \end{aligned}$$

from which we get

$$\begin{aligned} & \frac{\partial u}{\partial \tau} + \frac{\Delta\tau}{2} \frac{\partial^2 u}{\partial \tau^2} - \left( \frac{2r}{\sigma^2} + \left(\frac{\tilde{\sigma}}{\sigma}\right)^2 \right) \frac{\partial u}{\partial y} - \Delta\tau \left( \frac{2r}{\sigma^2} + \left(\frac{\tilde{\sigma}}{\sigma}\right)^2 \right) \frac{\partial^2 u}{\partial y \partial \tau} \\ & - \left(\frac{\tilde{\sigma}}{\sigma}\right)^2 \frac{\partial^2 u}{\partial y^2} + O[(\Delta\tau)^2, (\Delta y)^2] = 0 \end{aligned}$$

It is obvious that if  $\Delta\tau, \Delta y \rightarrow 0$ , then the original PDE (2) is recovered. Therefore, the Laasonen finite difference scheme, Equation (8), is consistent. Now according to Lax's equivalence theorem, [55], LFDS is convergent for all values of the parameters. Similar arguments hold for DFFDS and FVFIS. On the other hand, FVES and FVCNS are also convergent if the conditions (20) and (21) respectively, are satisfied.

TABLE 1. Call option prices using the Leland volatility model.

$S_0$	Exact (Linear)	Finite Difference Schemes		Finite Volume Schemes		
		DFFDS	LFDS	FVES	FVFIS	FVCNS
37.00	0.00001	0.04734	0.04893	0.00000	0.00054	0.00050
47.00	0.00182	0.30914	0.31368	0.00006	0.01340	0.01297
57.00	0.05078	1.09349	1.09949	0.00036	0.10771	0.10588
67.00	0.45226	2.93576	2.94472	0.00335	0.65191	0.64795
77.00	1.97686	6.06630	6.07104	0.01709	2.26632	2.26230
87.00	5.46222	10.56460	10.56947	1.25649	5.63768	5.63688
97.00	11.17037	16.34912	16.34757	6.49278	11.07370	11.07714
107.00	18.71972	23.33442	23.33541	16.15278	18.66210	18.66616
117.00	27.48006	31.19401	31.19406	26.33380	27.42028	27.42359
127.00	36.91158	39.65579	39.65486	36.72393	36.83405	36.83640
137.00	46.67034	48.63231	48.63285	46.68664	46.59745	46.59938
147.00	56.57397	57.89847	57.89955	56.61703	56.45705	56.45879
157.00	66.53723	67.45340	67.45432	66.59903	66.46595	66.46767
167.00	76.52370	77.08904	77.08984	76.50332	76.39766	76.39940
177.00	86.51886	86.88858	86.88931	86.48557	86.40846	86.41023
187.00	96.51716	96.75424	96.75478	96.48843	96.43330	96.43508
197.00	106.51657	106.59706	106.59728	106.40288	106.36023	106.36202
207.00	116.51636	116.67091	116.67078	116.55261	116.52319	116.52499
217.00	126.51629	126.48888	126.48906	126.39973	126.37558	126.37738
227.00	136.51627	136.46008	136.46065	136.39791	136.37870	136.38049
237.00	146.51626	146.57401	146.57489	146.53847	146.52419	146.52597
247.00	156.51626	156.41475	156.41484	156.38607	156.37369	156.37545
257.00	166.51626	166.40053	166.39979	166.37904	166.36868	166.37039
267.00	176.51626	176.52571	176.52411	176.51173	176.50347	176.50515

## 7. RESULTS AND DISCUSSIONS

In this section, we choose the same parameters:  $r = 0.1$ ,  $\sigma = 0.2$ ,  $K = 100$ ,  $T = 1$ ,  $\mu = 0.05$ ,  $\Delta t = 0.01$ ,  $a = 0.02$ ,  $M = 0.01$ , and  $C = 30$ , as illustrated in the literature [46]. Then we calculate the call option values using the proposed schemes, described in previous section 4, for different volatility models. We compare the approximate results with the exact value of the linear Black-Scholes model and among themselves also.

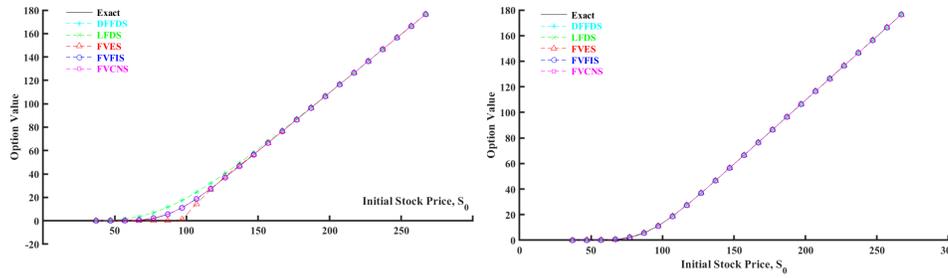


FIGURE 1. Approximate results of Equation (1) by using (a) Boyle and Vorst volatility model, and (b) Barles and Soner volatility model.

From Table 1 and Figure 8.7 (see Appendix), we observe that fully implicit FVS and Crank-Nicolson FVS provide comparatively better results than the other methods. Note that all of the methods provide poor results when the initial stock price is less than the strike price (here strike price, in comparison with the exact value of the linear Black-Scholes model).

From Table 8.3 in Appendix 8 and Figure 1 (a), we can make similar comments, but here the FVES gives a very poor approximation than the other methods when the initial stock price is less than the strike price ( $K = 100$ ). Table 8.4 in Appendix 8 and Figure 1(b) show that all of the methods provide a closer approximation to the exact value of the linear Black-Scholes model for all of the initial stock price, whether it is greater than the strike price,  $K = 100$ .

TABLE 2. Call option prices using RAPM volatility model.

$S_0$	$E_{\text{Exact}}$ (Linear)	Finite Difference Schemes		Finite Volume Schemes		
		DFFDS	LFDS	FVES	FVFIS	FVCNS
37.00	0.00001	0.09975	0.10198	0.00075	0.00054	0.00050
47.00	0.00182	0.64320	0.64859	0.02029	0.01340	0.01297
57.00	0.05078	2.01164	2.01581	0.16518	0.10771	0.10588
67.00	0.45226	4.58788	4.59820	0.97606	0.65191	0.64795
77.00	1.97686	8.35995	8.36458	3.28335	2.26632	2.26230
87.00	5.46222	13.24558	13.25523	7.65193	5.63768	5.63688
97.00	11.17037	19.14277	19.14296	13.90023	11.07370	11.07714
107.00	18.71972	25.95302	25.96004	21.59995	18.66210	18.66616
117.00	27.48006	33.49710	33.50138	30.05441	27.42028	27.42359
127.00	36.91158	41.58132	41.58229	39.02457	36.83405	36.83640
137.00	46.67034	50.16133	50.16454	48.32523	46.59745	46.59938
147.00	56.57397	59.07941	59.08280	57.80665	56.45705	56.45879
157.00	66.53723	68.31796	68.31997	67.49623	66.46595	66.46767
167.00	76.52370	77.71942	77.72089	77.20105	76.39767	76.39941
177.00	86.51886	87.32671	87.32812	87.03280	86.40846	86.41023
187.00	96.51716	97.04281	97.04399	96.91429	96.43330	96.43509
197.00	106.51657	106.79810	106.79881	106.75046	106.36023	106.36202
207.00	116.51636	116.77932	116.77954	116.81690	116.52319	116.52498
217.00	126.51629	126.56491	126.56478	126.62193	126.37555	126.37734
227.00	136.51627	136.50685	136.50637	136.57895	136.37864	136.38041
237.00	146.51626	146.59198	146.59115	146.67807	146.52410	146.52585
247.00	156.51626	156.42589	156.42489	156.50633	156.37350	156.37521
257.00	166.51626	166.40454	166.40336	166.47896	166.36838	166.37004
267.00	176.51626	176.52231	176.52094	176.59035	176.50307	176.50468

Finally, from Table 2 and the corresponding Figure 8.8 in Appendix 8, we may observe that for the RAPM volatility model, the FVCNS and FVFIS give better approximation than the other numerical schemes when the initial stock price is closer to and/or greater than the strike price. On the other hand, from Figures 2, 3, 4, it is clear that FVFIS and FVCNS produce comparatively better results than the other schemes for all of the volatility models. Figures 5, 6 depict the option prices at various time periods (from initial time  $t = 0$  to maturity time,  $t = T$ ) with different initial stock values. The similar results of solution surface for option price by using Barles and Soner volatility model and RAPM volatility model are presented in Appendix 8, see Figures 8.9,8.10.

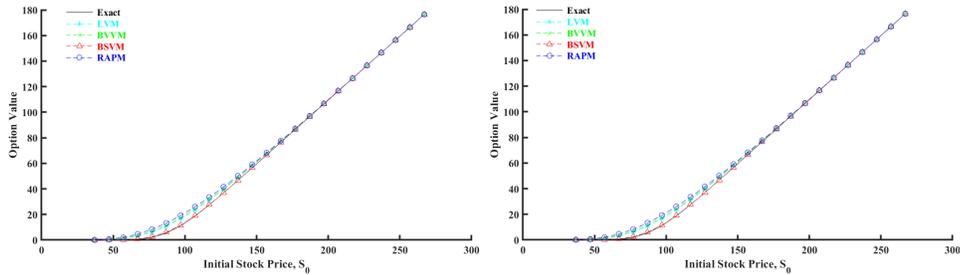


FIGURE 2. Approximate results of Equation (1) using (a) Dufort-Frankel Finite Difference Scheme, and (b) Laasonen Finite Difference Scheme.

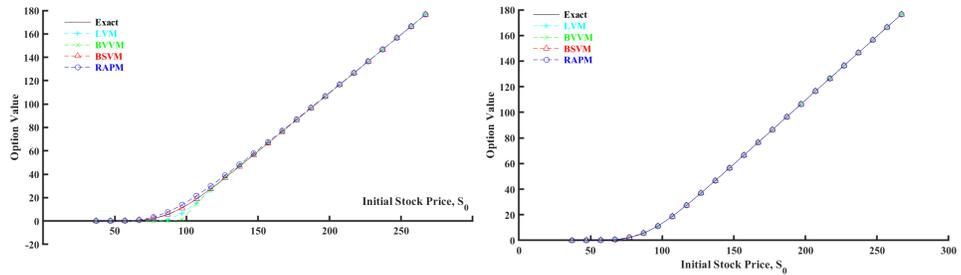


FIGURE 3. Approximate results of Equation (1) using (a) Finite Volume Explicit Scheme, and (b) Finite Volume Fully Implicit Scheme.

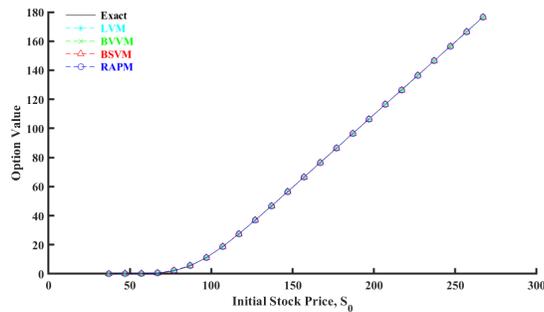


FIGURE 4. Approximate results of Equation (1) using Finite Volume Crank-Nicolson Scheme.

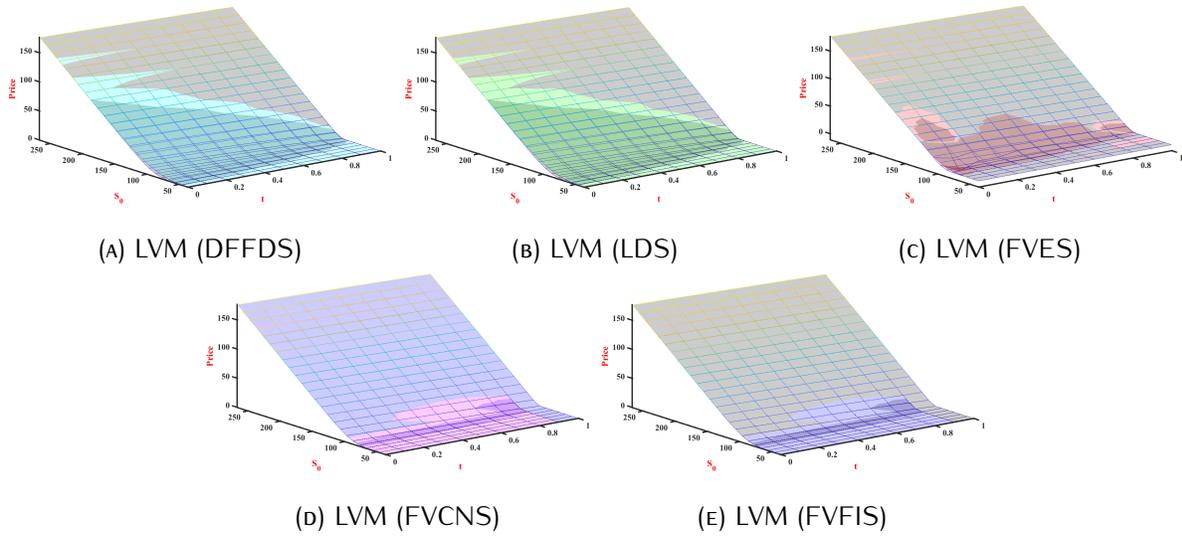


FIGURE 5. Solution surface for option price by using Leland volatility model.

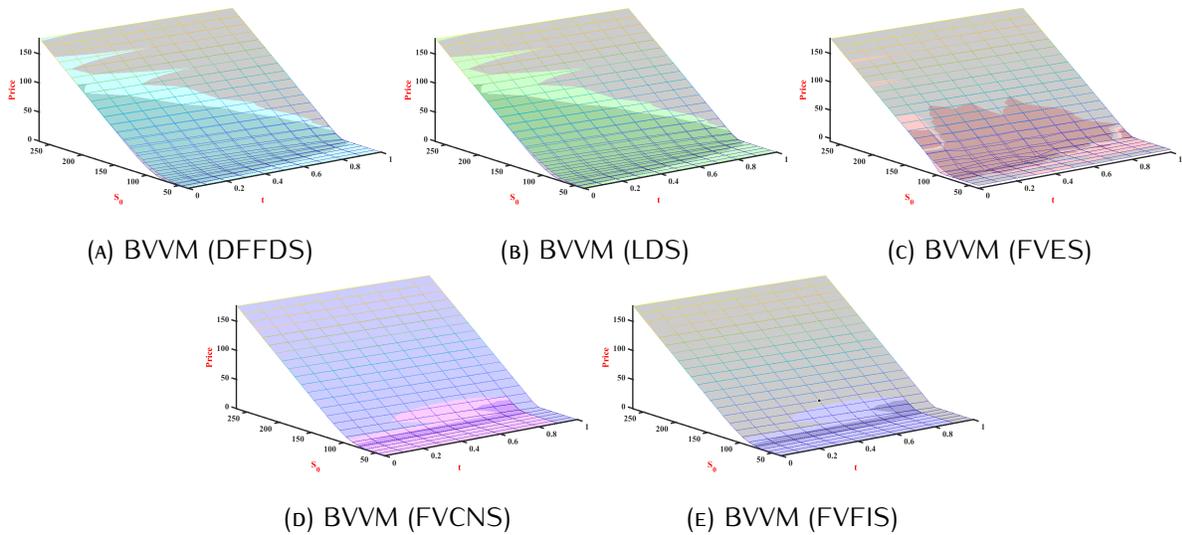


FIGURE 6. Solution surface for option price by using Boyle and Vorst volatility model.

### 8. CONCLUSION

In this research work, we have derived some numerical schemes using the FVM and FDM to solve the non-linear Black-Scholes PDE for European option pricing with the transaction costs by exploiting the transformations available in the existing literature [46]. Thus we have modified the model equation accordingly to a non-linear parabolic PDE. For the convergence of these schemes, stability and consistency have been shown rigorously. Then these schemes have been applied to various volatility models. According to the visible results, as presented in the earlier sections, it

is noted that all of the proposed schemes provide the best approximation to the exact value of the linear Black-Scholes model for all initial stock prices, regardless of whether they are closer to or greater than the strike price; particularly in the case of Barles and Soner Volatility Model. We may claim that the FVFIS and FVCNS approximate better than the other methods for all four-volatility models. Thus, it is observed that the FVFIS and FVCNS are very effective and proficient in locating approximate solutions to non-linear Black-Scholes models. Notice that the limitation of these schemes is that they may offer poor results sometimes when the initial stock price is less than the strike price compared to the exact value of the linear Black-Scholes model. Finally, we may conclude that the proposed schemes may be applied to other non-linear partial differential equations to compute the numerical solutions with the desired accuracy.

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APPENDIX

This section contains the supporting figures and tables to observe the accuracy of the solution methodologies.

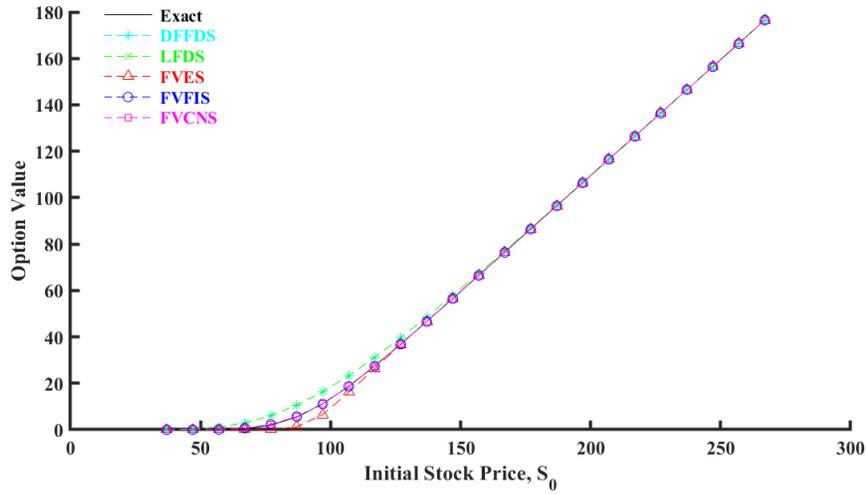


FIGURE 8.7. Approximate results of Equation (1) by using Leland volatility model.

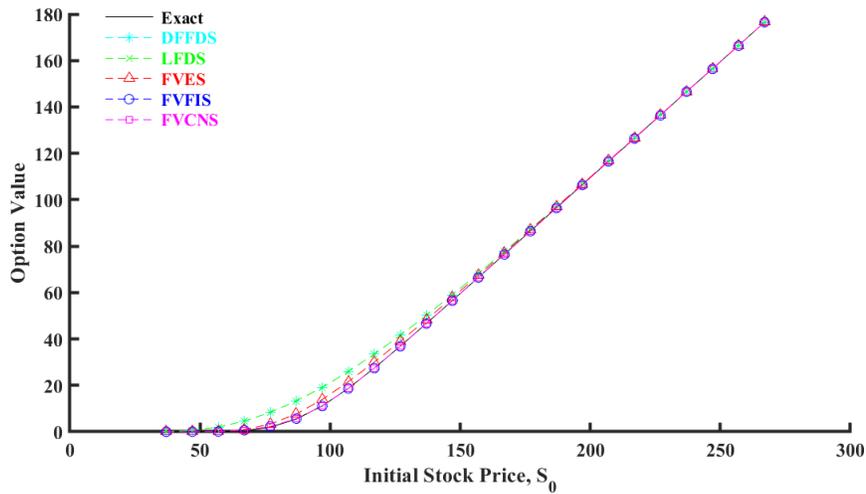


FIGURE 8.8. Approximate results of Equation (1) by using RAPM volatility model.

TABLE 8.3. Call option prices using Boyle and Vorst volatility model.

$S_0$	Exact (Linear)	Finite Difference Schemes		Finite Volume Schemes		
		DFFDs	LFDS	FVES	FVFIS	FVCNS
37.00	0.00001	0.08130	0.08362	0.00000	0.00054	0.00050
47.00	0.00182	0.44919	0.45475	0.00016	0.01340	0.01297
57.00	0.05078	1.43187	1.43827	0.00137	0.10771	0.10588
67.00	0.45226	3.52436	3.53389	0.00181	0.65191	0.64795
77.00	1.97686	6.88032	6.88512	-0.01095	2.26632	2.26230
87.00	5.46222	11.52515	11.53065	-0.03314	5.63768	5.63688
97.00	11.17037	17.36380	17.36252	1.18226	11.07370	11.07714
107.00	18.71972	24.30605	24.30805	14.49730	18.66210	18.66616
117.00	27.48006	32.07120	32.07189	26.66264	27.42028	27.42359
127.00	36.91158	40.41514	40.41449	37.98090	36.83405	36.83640
137.00	46.67034	49.26375	49.26507	47.25031	46.59745	46.59938
147.00	56.57397	58.41457	58.41649	56.80881	56.45705	56.45879
157.00	66.53723	67.86188	67.86337	66.66025	66.46595	66.46767
167.00	76.52370	77.41353	77.41510	76.50821	76.39767	76.39940
177.00	86.51886	87.14056	87.14251	86.47640	86.40846	86.41023
187.00	96.51716	96.94695	96.94884	96.47214	96.43330	96.43509
197.00	106.51657	106.75019	106.75144	106.38908	106.36023	106.36202
207.00	116.51636	116.78198	116.78255	116.54146	116.52319	116.52500
217.00	126.51629	126.57938	126.58055	126.39087	126.37558	126.37738
227.00	136.51627	136.53059	136.53250	136.39143	136.37869	136.38048
237.00	146.51626	146.62463	146.62715	146.53436	146.52418	146.52596
247.00	156.51626	156.45839	156.45983	156.38285	156.37367	156.37541
257.00	166.51626	166.43684	166.43715	166.37675	166.36863	166.37034
267.00	176.51626	176.55434	176.55347	176.51042	176.50341	176.50508

TABLE 8.4. Call option prices using Barles and Soner volatility model.

$S_0$	Exact (Linear)	Finite Difference Schemes		Finite Volume Schemes		
		DFFDs	LFDS	FVES	FVFIS	FVCNS
37.00	0.00001	0.00056	0.00060	0.00047	0.00054	0.00050
47.00	0.00182	0.01448	0.01496	0.01255	0.01340	0.01297
57.00	0.05078	0.11766	0.11964	0.10402	0.10771	0.10588
67.00	0.45226	0.70853	0.71289	0.64387	0.65191	0.64795
77.00	1.97686	2.42414	2.42846	2.25831	2.26632	2.26230
87.00	5.46222	5.90671	5.90840	5.63756	5.63768	5.63688
97.00	11.17037	11.38951	11.38681	11.08596	11.07370	11.07714
107.00	18.71972	18.90018	18.89694	18.68219	18.66210	18.66616
117.00	27.48006	27.57177	27.56903	27.44285	27.42028	27.42359
127.00	36.91158	36.90926	36.90718	36.85630	36.83405	36.83640
137.00	46.67034	46.63600	46.63420	46.61611	46.59745	46.59938
147.00	56.57397	56.47483	56.47316	56.47209	56.45705	56.45879
157.00	66.53723	66.47486	66.47318	66.47729	66.46595	66.46767
167.00	76.52370	76.40243	76.40071	76.40662	76.39766	76.39940
177.00	86.51886	86.41173	86.40997	86.41573	86.40846	86.41022
187.00	96.51716	96.43561	96.43381	96.43938	96.43329	96.43508
197.00	106.51657	106.36234	106.36053	106.36582	106.36022	106.36202
207.00	116.51636	116.52509	116.52327	116.52827	116.52319	116.52499
217.00	126.51629	126.37745	126.37563	126.38047	126.37557	126.37737
227.00	136.51627	136.38055	136.37873	136.38340	136.37869	136.38048
237.00	146.51626	146.52603	146.52421	146.52870	146.52418	146.52596
247.00	156.51626	156.37561	156.37380	156.37798	156.37368	156.37543
257.00	166.51626	166.37066	166.36886	166.37271	166.36866	166.37038
267.00	176.51626	176.50552	176.50374	176.50725	176.50346	176.50514

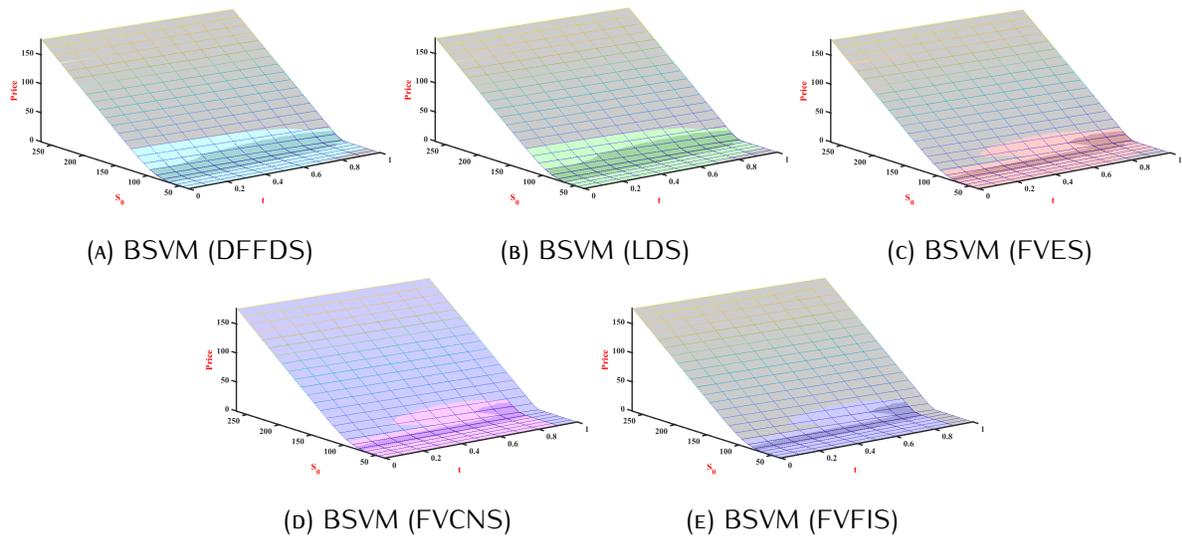


FIGURE 8.g. Solution surface for option price by using Barles and Soner volatility model.

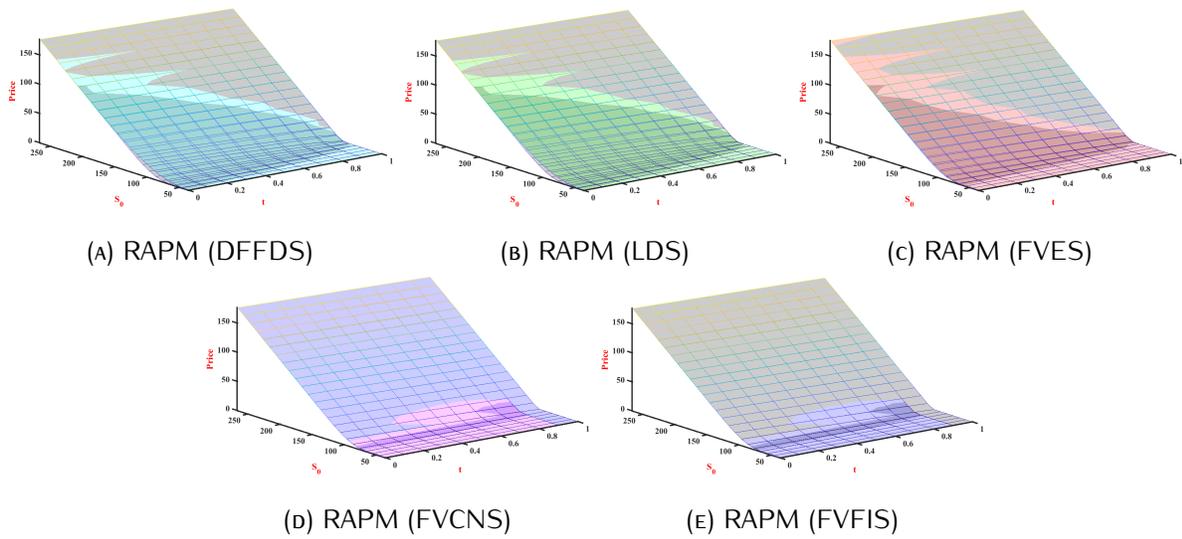


FIGURE 8.10. Solution surface for option price by using RAPM volatility model.